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1989 J. Phys. A: Math. Gen. 22 L255

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LETTER TO THE EDITOR

Solving formula of the two-dimensional Toda lattice

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Received 15 December 1988

Abstract. A solving formula of the two-dimensional Toda lattice is proved, making use of the direct method. This formula is a generalisation and unification of the Bäcklund transformation, the non-linear superposition formula and the Casorati determinant solutions.

It is known that the two-dimensional Toda lattice has Casorati determinant solutions which are generalisations of the usual N -soliton solutions [1]. A linear Bäcklund transformation was found which generates an $(N + 1)$ -soliton solution from an N -soliton solution [2, 3]. In this letter, we shall give a generalisation and unification of the above results, making use of the direct substitution means using Hirota's bilinear form [4].

We have the following bilinear form for the two-dimensional Toda lattice

$$f_n(\partial_x \partial_s f_n) - (\partial_x f_n)(\partial_s f_n) - (f_{n+1}f_{n-1} - f_n^2) = 0. \tag{1}$$

Now, given that g_n is a solution of (1), f_n^m ($m = 0, 1, \dots$) satisfy

$$\partial_x f_n^m = \frac{g_n}{g_{n+1}} f_{n+1}^m + \frac{\partial_x g_{n+1}}{g_{n+1}} f_n^m \tag{2}$$

and

$$\partial_s f_n^m = -\frac{g_{n+1}}{g_n} f_{n-1}^m + \frac{\partial_s g_n}{g_n} f_n^m. \tag{3}$$

We define

$$r_n^m = \begin{vmatrix} f_n^0 & f_{n+1}^0 & \dots & f_{n+m}^0 \\ f_n^1 & f_{n+1}^1 & \dots & f_{n+m}^1 \\ \vdots & \vdots & \dots & \vdots \\ f_n^m & f_{n+1}^m & \dots & f_{n+m}^m \end{vmatrix} (g_{n+1} \dots g_{n+m})^{-1} \\ = |n, n+1, \dots, n+m| (g_{n+1} \dots g_{n+m})^{-1} \tag{4}$$

where the notation $|n_1, n_2, \dots, n_k|$ is as used in [4]. Using (2) and (3), we have

$$\partial_x r_n^m = \frac{|n, \dots, n+m-1, n+m+1|}{g_{n+1} \dots g_{n+m-1} g_{n+m+1}} + \frac{(\partial_x g_{n+m+1})|n, \dots, n+m|}{g_{n+1} \dots g_{n+m+1}} \tag{5}$$

$$\partial_s r_n^m = -\frac{|n-1, n+1, \dots, n+m|}{g_n g_{n+2} \dots g_{n+m}} + \frac{(\partial_s g_n)|n, \dots, n+m|}{g_n \dots g_{n+m}} \tag{6}$$

$$\begin{aligned} \partial_x \partial_s r_n^m = & \frac{1}{g_n \dots g_{n+m}} \left(\frac{g_n (\partial_x \partial_s g_n) - (\partial_x g_n) (\partial_s g_n) - g_{n+1} g_{n-1}}{g_n} |n, \dots, n+m| \right. \\ & + \frac{g_{n+1} g_{n+m}}{g_{n+m+1}} |n-1, n+1, \dots, n+m-1, n+m+1| \\ & - \frac{g_{n+1} (\partial_x g_{n+m+1})}{g_{n+m+1}} |n-1, n+1, \dots, n+m| \\ & \left. + \frac{(\partial_s g_n) (\partial_x g_{n+m+1})}{g_{n+m+1}} |n, \dots, n+m| \right) \end{aligned} \tag{7}$$

$$\begin{aligned} (\partial_x r_n^m) (\partial_s r_n^m) = & \frac{|n, \dots, n+m|}{g_n g_{n+1}^2 \dots g_{n+m}^2 g_{n+m+1}} [|n, \dots, n+m-1, n+m+1| g_{n+m} (\partial_s g_n) \\ & + (\partial_s g_n) (\partial_x g_{n+m+1}) |n, \dots, n+m| \\ & - g_{n+1} (\partial_x g_{n+m+1}) |n-1, n+1, \dots, n+m|] \\ & - \frac{|n-1, n+1, \dots, n+m| |n, \dots, n+m-1, n+m+1|}{g_n \dots g_{n+m-1} g_{n+2} \dots g_{n+m+1}} \end{aligned} \tag{8}$$

and

$$\begin{aligned} r_n^m (\partial_x \partial_s r_n^m) - (\partial_x r_n^m) (\partial_s r_n^m) - r_{n+1}^m r_{n-1}^m + r_n^m r_n^m \\ = \left(\frac{|n, \dots, n+m|}{g_n \dots g_{n+m}} \right)^2 [g_n (\partial_x \partial_s g_n) - (\partial_x g_n) (\partial_s g_n) - g_{n+1} g_{n-1} + g_n^2] \\ + \frac{1}{g_{n+2} \dots g_{n+m+1}} \frac{1}{g_n \dots g_{n+m-1}} \\ \times (|n, \dots, n+m| |n-1, n+1, \dots, n+m-1, n+m+1| \\ - |n-1, n+1, \dots, n+m| |n, \dots, n+m-1, n+m+1| \\ - |n+1, \dots, n+m+1| |n-1, \dots, n+m-1|) \\ = - \frac{1}{g_{n+2} \dots g_{n+m+1}} \frac{1}{g_n \dots g_{n+m-1}} \\ \times \left| \begin{array}{ccc|ccc} n-1 & | & n & | & n+1, \dots, n+m-1 & | \\ n-1 & | & n & | & 0 & | \\ \hline & & & & n+1, \dots, n+m-1 & | \\ n+m & | & n+m+1 & | & & | \\ \hline & & & & n+m & | \\ & & & & n+m+1 & | \end{array} \right| \\ = 0. \end{aligned} \tag{9}$$

Thus r_n^m ($m=0, 1, \dots$) are solutions of (1).

If we take $g_n = 1$, then we get the Casorati determinant solution [1]

$$r_n^m = \begin{vmatrix} f_n^0 & \dots & f_{n+m}^0 \\ \vdots & \ddots & \vdots \\ f_n^m & \dots & f_{n+m}^m \end{vmatrix}$$

where $\partial_x f_n^m = f_{n+1}^m$ and $\partial_s f_n^m = -f_{n-1}^m$.

If we take $m = 0$, then we know f_n are solutions of (1) [2].

If we take $m = 1$, then we get the non-linear superposition formula [5]

$$g_n f_n^1 = \exp(D_n) f_n^1 f_n^0$$

where D_n is Hirota's bilinear difference operator.

The author would like to express his sincere thanks to Professor Ben-Yu Guo for encouragement and advice.

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